

Entropic Uncertainty Relations and their Applications in Quantum Cryptography

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Take Home Message

Entropic uncertainty relations allow us to formalize physical intuition in information theoretic terms.

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The min-entropy is the preferred measure of uncertainty for cryptographic applications.

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Uncertainty relations in terms of the min-entropy consequently have a wide range of applications in quantum cryptography.

① Entropic Uncertainty Relations

A Simple Model

Uncertainty for Classical Observers

Maassen & Uffink Relation

Rényi Entropies and the Min-Entropy

② Uncertainty for Quantum Observers

Uncertainty vs. Entanglement

Uncertainty Relation for the Min-Entropy

Quantum Rényi Entropies

Generalized Maassen & Uffink Relation

③ Application to Quantum Key Distribution

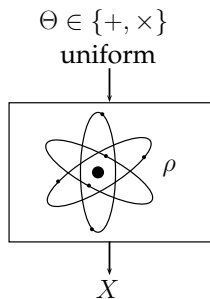
Setup for Quantum Key Distribution

Security Proof Sketch

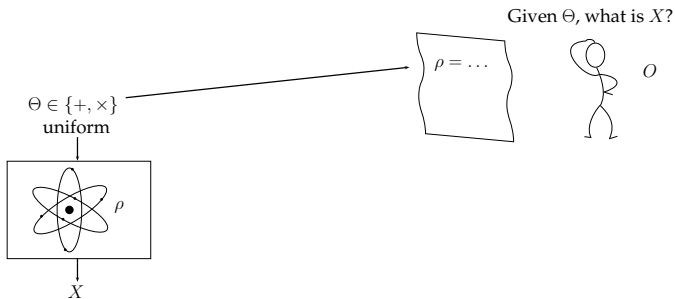
④ More Entropic Uncertainty Relations

From Quantum Systems to Random Variables

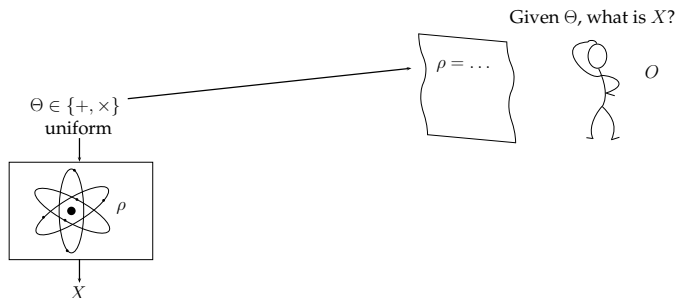
- Simplest example: Measure photon polarization in rectilinear '+' or diagonal '×' basis.
- More generally, we consider two arbitrary measurements on an arbitrary quantum system in state ρ .



Uncertainty Relation with Classical Observer

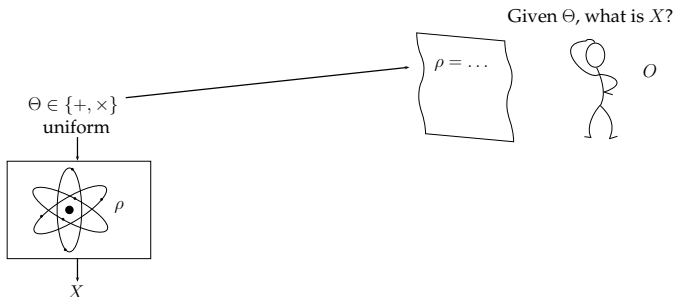


Uncertainty Relation with Classical Observer



- Heisenberg's uncertainty principle [Hei27]:
 It is impossible to predict the outcome of both measurements.

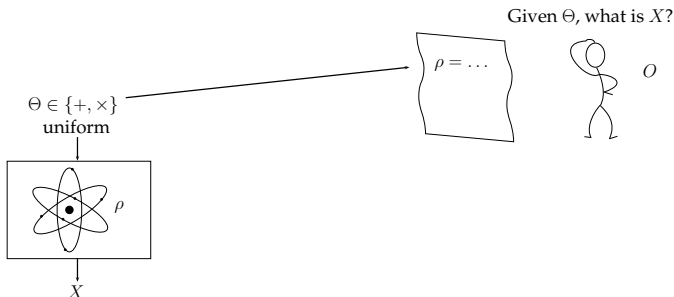
Uncertainty Relation with Classical Observer



- Heisenberg's uncertainty principle [Hei27]:
It is impossible to predict the outcome of both measurements.
- Deutsch [Deu83] proposed to characterize this uncertainty in terms of the Shannon entropy [Sha48].

$$H(X) := \sum_x P_X(x) \log \frac{1}{P_X(x)}, \quad H(X|Y) := \sum_y P_Y(y) H(X|Y=y).$$

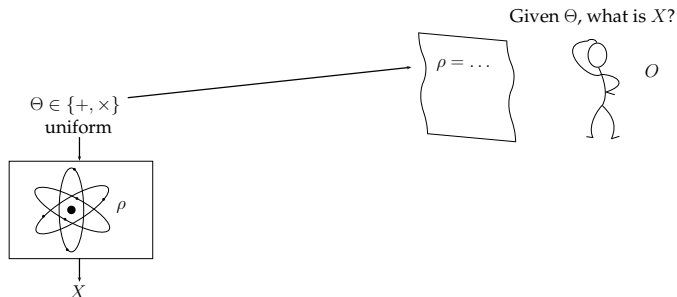
Uncertainty Relation with Classical Observer



- Simplest example: measure polarization for photon prepared in the rectilinear '+' basis.

$$H(X|O, \Theta) = \frac{1}{2}H(X|O, \Theta = '+') + \frac{1}{2}H(X|O, \Theta = '\times').$$

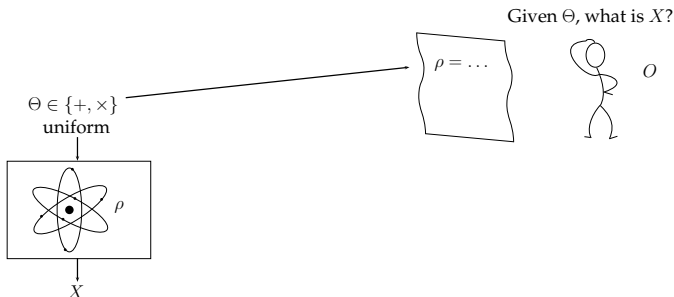
Uncertainty Relation with Classical Observer



- Simplest example: measure polarization for photon prepared in the rectilinear '+' basis.

$$H(X|O, \Theta) = \frac{1}{2} \underbrace{H(X|O, \Theta = '+')}_{= 0} + \frac{1}{2} \underbrace{H(X|O, \Theta = '\times')}_{= 1}.$$

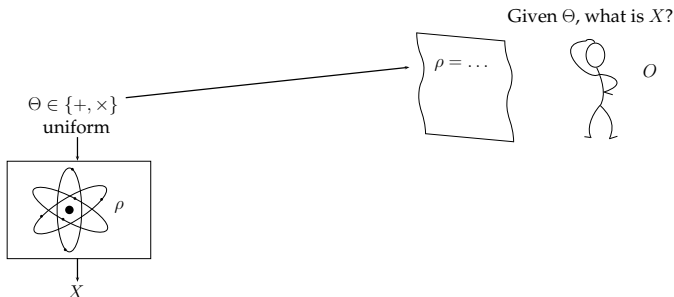
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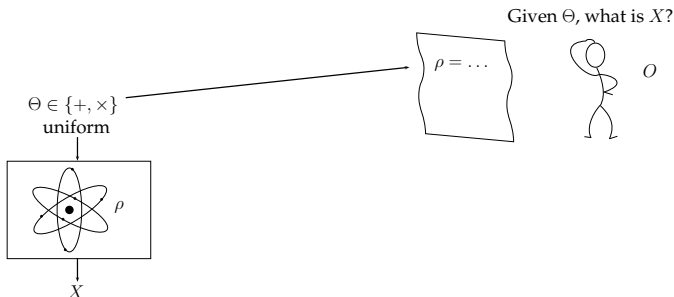


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$$H(X|O, \Theta) = \frac{1}{2}H(X|O, \Theta = '+') + \frac{1}{2}H(X|O, \Theta = '\times') = \frac{1}{2}.$$

- In fact, the bound $H(X|O, \Theta) \geq \frac{1}{2}$ holds independently of how the state ρ is prepared!

Uncertainty Relation with Classical Observer



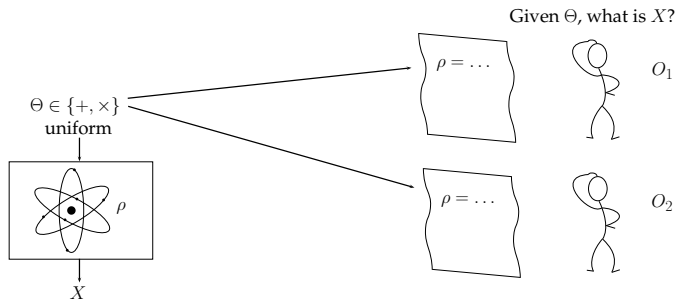
More generally, for arbitrary projective measurements:

Maassen & Uffink [MU88]

$$H(X|O, \Theta) \geq \frac{1}{2} \log \frac{1}{c}, \quad c = \max_{x,y} |\langle \phi_x^+ | \phi_y^x \rangle|^2,$$

where c is called *overlap* and $\{|\phi_x^+\rangle\}_x$ and $\{|\phi_y^x\rangle\}_y$ are the eigenvectors of the respective measurements.

Uncertainty Relation with Two Classical Observers

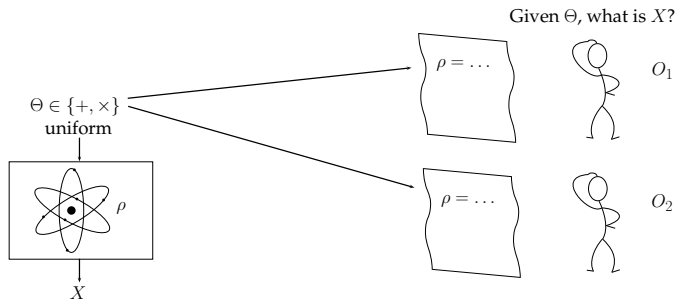


The previous relations trivially extends to two observers:

Maassen & Uffink [MU88]

$$H(X|O_1, \Theta) + H(X|O_2, \Theta) \geq \log_2 \frac{1}{c}.$$

Uncertainty Relation with Two Classical Observers



More generally, the relation also holds for Rényi entropies [Rén61]:

Maassen & Uffink [MU88]

$$H_\alpha(X|O_1, \Theta) + H_\beta(X|O_2, \Theta) \geq \log \frac{1}{c}, \quad \text{where} \quad \frac{1}{\alpha} + \frac{1}{\beta} = 2.$$

Classical Rényi Entropy Intermezzo

- Rényi [Rén61] defined a family of divergences for $\alpha \in [0, \infty]$ as

$$D_\alpha(P\|Q) := \frac{1}{\alpha - 1} \log \sum_z P(z)^\alpha Q(z)^{1-\alpha},$$

taking limits where necessary.

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- The Rényi entropy is related to the divergence,

$$H_\alpha(X) := \frac{1}{1 - \alpha} \log \sum_x P_X(x)^\alpha = -D_\alpha(P_X\|1_X).$$

- The proper extension to conditional Rényi entropy is disputed. Here we employ a definition due to Arimoto [Ari73]:

$$H_\alpha(X|Y) := \frac{\alpha}{1 - \alpha} \log \left(\sum_y P_Y(y) \exp \left(\frac{1 - \alpha}{\alpha} H_\alpha(X|Y=y) \right) \right)$$

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- This extension is closely related to the Rényi divergence,

$$H_\alpha(X|Y) = \sup_{Q_Y} -D_\alpha(P_{XY}\|1_X \times Q_Y).$$

Classical Rényi Entropy Intermezzo

- The function $\alpha \rightarrow H_\alpha(X|Y)$ is monotonically decreasing.
- The Shannon entropy is $\lim_{\alpha \rightarrow 1} H_\alpha(X|Y) = H(X|Y)$.
- It satisfies a data-processing inequality. For any channel $Y \rightarrow Z$, we have

$$H_\alpha(X|Y) \leq H_\alpha(X|Z).$$

Operationally, the uncertainty of X increases when we manipulate the side information Y .

- The quantity has found operational significance. In particular,

$$\lim_{\alpha \rightarrow \infty} H_\alpha(X|Y) = H_{\min}(X|Y) = -\log p_{\text{guess}}(X|Y),$$

where $p_{\text{guess}}(X|Y) = \sum_y P_Y(y) \max_x P_{X|Y}(x)$ is the (average) probability of correctly guessing X given Y .

Min-Entropy and “Leftover Hashing”

- The omnipresence of the min-entropy in cryptography is due to the “Leftover Hashing” Lemma.

Bennett *et al.* [BBR88] & Impagliazzo *et al.* [ILL89, IZ89]

Let $X \in \mathcal{X}$ be such that $H_{\min}(X|E) \geq k$, and $\ell \in \mathbb{N}$. Then, there exists a family of Hash function $\{f_S\}_S$ from \mathcal{X} to $\{0, 1\}^\ell$ such that $f_S(X)$ is δ -close to uniform and independent of E and S , where

$$\delta \sim \sqrt{2^{\ell-k}}$$

and S is chosen uniformly at random.

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$$\delta \sim \sqrt{2^{\ell-k}} + 2\varepsilon$$

and S is chosen uniformly at random.

- This is “tight” in the sense that one cannot expect to extract more than $H_{\min}^{\varepsilon}(X|E)$ bits of uniform and independent randomness.

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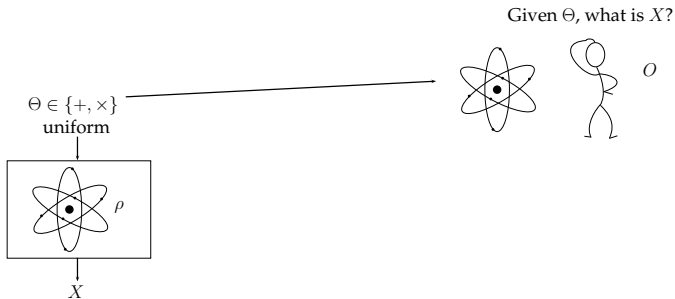
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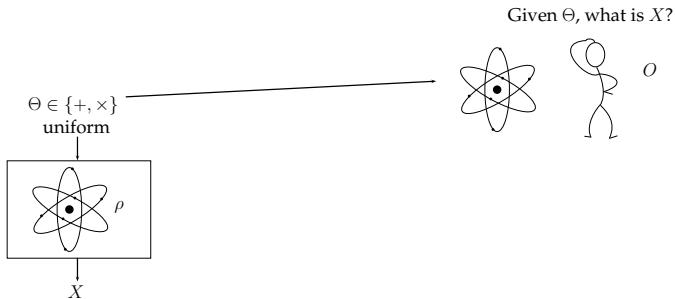
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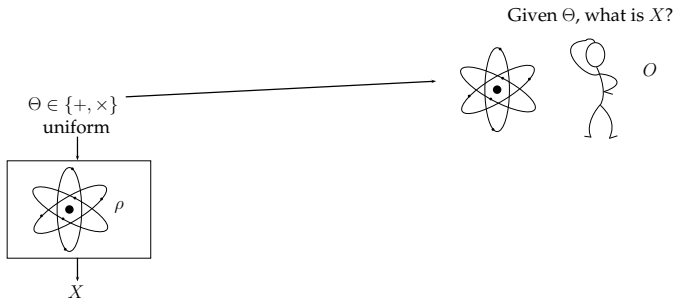
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Uncertainty vs. Entanglement



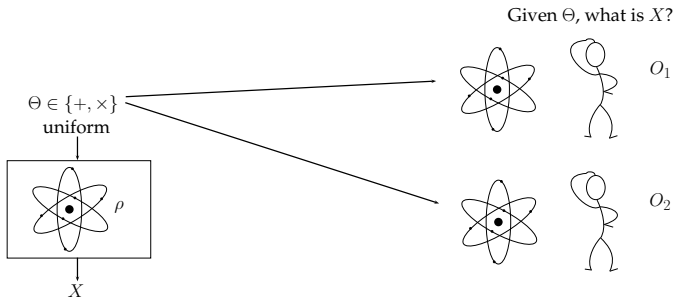
- Consider a joint quantum state ρ_{AO} .
- If $\rho_{AO} = |\psi\rangle\langle\psi|$ is maximally entangled then $H(X|O, \Theta) = 0$.
 - Why? The observer may choose a measurement on O — depending on Θ — to get perfect correlation with X .

Uncertainty vs. Entanglement



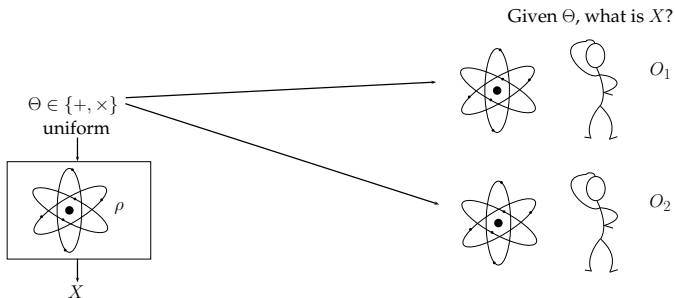
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 - Why? The observer may choose a measurement on O — depending on Θ — to get perfect correlation with X .
- No uncertainty relation here!

Uncertainty and Monogamy of Entanglement



Can monogamy of entanglement help?

Uncertainty and Monogamy of Entanglement



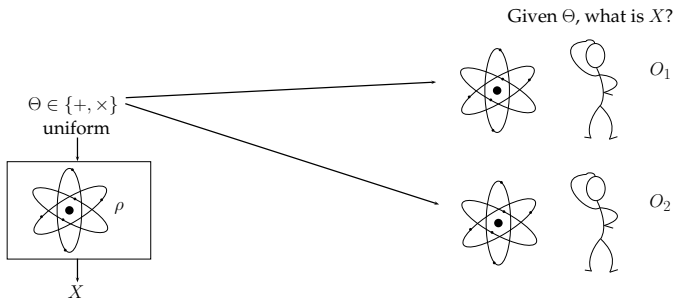
Can monogamy of entanglement help?

Berta, Christandl, Colbeck, Renes & Renner [BCC⁺10]

The post-measurement states satisfy

$$H(X|O_1, \Theta) + H(X|O_2, \Theta) \geq \log_2 \frac{1}{c}.$$

Uncertainty and Monogamy of Entanglement



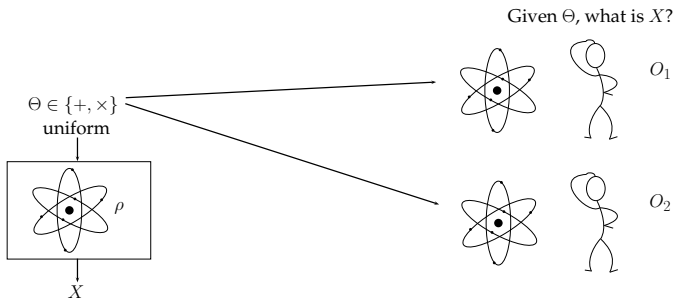
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$$H(X|O_1, \Theta) + H(X|O_2, \Theta) \geq \log_2 \frac{1}{c}.$$

- $H(A|B) := H(AB) - H(B)$ is the von Neumann entropy.
- The overlap c is defined as in the classical relation.

Uncertainty and Monogamy of Entanglement



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The post-measurement states satisfy

$$H(X|O_1, \Theta) + H(X|O_2, \Theta) \geq \log_2 \frac{1}{c}.$$

- Can this be generalized to quantum Rényi entropies, in particular the min-entropy?

Quantum Guessing Probability and the Min-Entropy

- The guessing probability naturally generalizes to the case of classical quantum states $\rho_{XE} = \bigoplus_x P(x)\rho_E^x$:

$$p_{\text{guess}}(X|E) := \sup \left\{ \sum_x P(x) \text{tr}(\rho_E^x M_E^x) \mid \{M_E^x\} \text{ a POVM} \right\}.$$

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- Then, we retain $H_{\min}(X|E) := -\log p_{\text{guess}}(X|E)$ as proposed by König, Renner & Schaffner [KRS09].
- The operational interpretation remains intact:

Renner [Ren05]: Quantum Leftover Hash Lemma

Let $X \in \mathcal{X}$ be such that $H_{\min}^\epsilon(X|E) \geq k$, and $\ell \in \mathbb{N}$. Then, there exists a family of Hash function $\{f_s\}_s$ from \mathcal{X} to $\{0, 1\}^\ell$ such that $f_S(X)$ is δ -close to uniform and independent of E and S , where

$$\delta \sim \sqrt{2^{\ell-k}} + 2\epsilon.$$

Uncertainty Relation for min-/max Entropies

T & Renner [TR11]: Min-Max Entropic Uncertainty

For any state $\rho_{AO_1O_2}$ and POVMs $\{M_x\}$ and $\{N_x\}$ on A:

$$H_{\min}(X|O_1, \Theta) + H_{\max}(X|O_2, \Theta) \geq \log_2 \frac{1}{c},$$

$$c = \max_{x,y} \left\| \sqrt{M_x} \sqrt{N_y} \right\|_{\infty}^2.$$

- Reduces to overlap defined previously for projective measurements.

Uncertainty Relation for min-/max Entropies

T & Renner [TR11]: Smooth Entropic Uncertainty

For any state $\rho_{AO_1O_2}$, $\varepsilon \geq 0$, and POVMs $\{M_x\}$ and $\{N_x\}$ on A:

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- Reduces to overlap defined previously for projective measurements.
- This also generalizes the previous result for the von Neumann entropy due to asymptotic equipartition [TCR09]:

$$\frac{1}{n} H_{\min}^{\varepsilon}(A^n|B^n), \frac{1}{n} H_{\max}^{\varepsilon}(A^n|B^n) \xrightarrow{n \rightarrow \infty} H(A|B).$$

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- What about general Rényi entropies?

Quantum Rényi Entropy Intermezzo

We want to find a quantum generalization of the Rényi entropy that satisfies the following

- It evaluates to the von Neumann, min-, and max-entropies for $\alpha = 1, \infty, \frac{1}{2}$, respectively.
- It has mathematical properties that we expect from an operational quantity, e.g. a data-processing inequality:

$$H_\alpha(A|B) \leq H_\alpha(A|C) \quad \text{for any quantum channel } \mathcal{E} : B \rightarrow C.$$

Such a generalization was recently proposed by Müller-Lennert, Dupuis, Szehr, Fehr & T [MLDS⁺13] (see also [WWY13]).

Uncertainty Relation for Quantum Rényi Entropy

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Müller-Lennert *et al.* [MLDS⁺13]: Quantum Rényi Entropy

The quantum conditional Rényi entropy is, for $\alpha \geq \frac{1}{2}$,

$$H_\alpha(A|B) := \sup_{\sigma_B} -D_\alpha(\rho_{AB} \| 1_A \otimes \sigma_B),$$

where $D_\alpha(\rho \| \sigma) := \frac{2\alpha}{\alpha-1} \log \left\| \rho^{\frac{1}{2}} \sigma^{\frac{1-\alpha}{2\alpha}} \right\|_{2\alpha}$ uses the Schatten norm.

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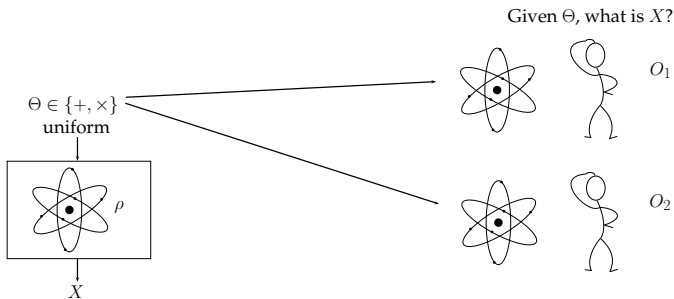
where $D_\alpha(\rho \| \sigma) := \frac{2\alpha}{\alpha-1} \log \left\| \rho^{\frac{1}{2}} \sigma^{\frac{1-\alpha}{2\alpha}} \right\|_{2\alpha}$ uses the Schatten norm.

- We recover the well-known special cases for $\alpha = \infty, \frac{1}{2}$:

$$H_{\min}(A|B) = \sup \left\{ \lambda \in \mathbb{R} \mid \exists \sigma_B : \rho_{AB} \leq 2^{-\lambda} 1_A \otimes \sigma_B \right\},$$

$$H_{\max}(A|B) = \sup_{\sigma_B} \left(2 \log F(\rho_{AB}, 1_A \otimes \sigma_B) \right).$$

Generalized Masseen & Uffink Relation



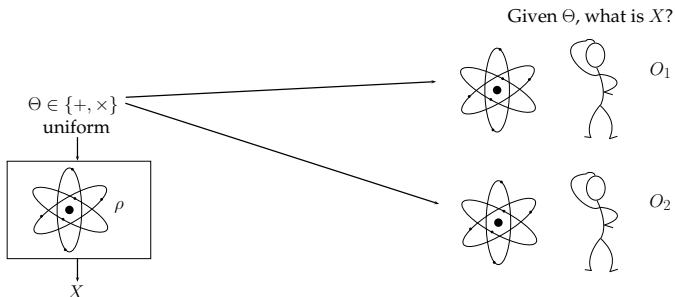
Müller-Lennert *et al.* [MLDS⁺13]: General Entropic Uncertainty

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$$\frac{1}{\alpha} + \frac{1}{\beta} = 2, \quad \text{and} \quad c = \max_{x,y} \left\| \sqrt{M_x} \sqrt{N_y} \right\|_\infty^2.$$

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$$H_\alpha(X|O_1, \Theta) + H_\beta(X|O_2, \Theta) \geq \log_2 \frac{1}{c}.$$

- Proof only uses a few basic properties of H_α and employs a strategy by Coles, Colbeck, Yu & Zwolak [CCYZ12].

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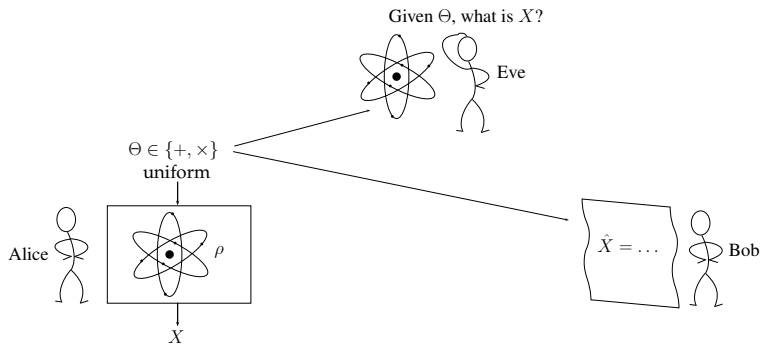
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Setup for Quantum Key Distribution

- We consider the entanglement-based Bennett-Brassard 1984 protocol [BB84, BBM92]

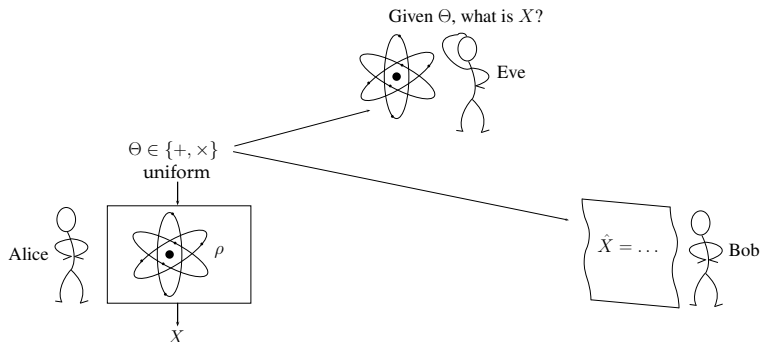
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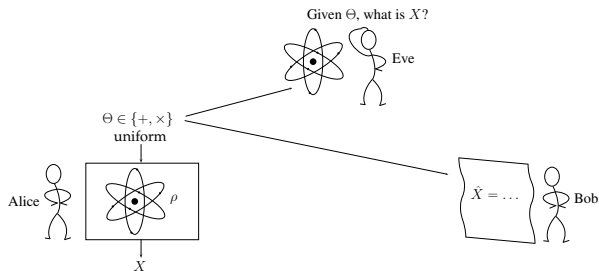
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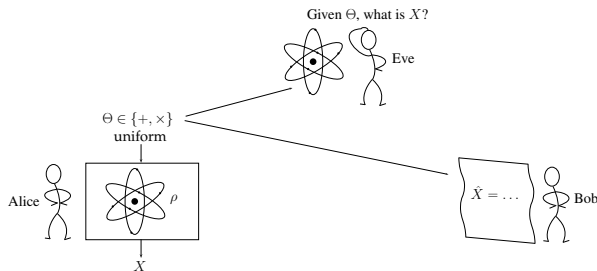
- Here, Θ , X and \hat{X} are n -bit strings.

Uncertainty Relation for n Systems



- Measure either all n systems in '+' or in 'x'.

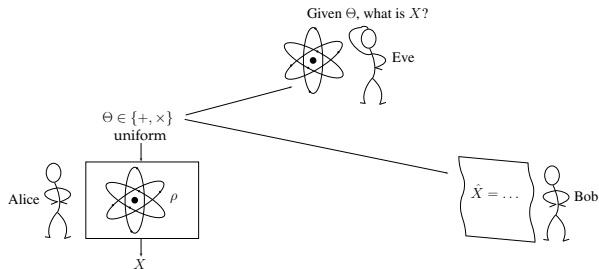
Uncertainty Relation for n Systems



- Measure either all n systems in '+' or in '×'.
- Define $F_0 := \{ '++ \dots +' , ' \times \times \dots \times ' \}$.

$$H_{\min}(X|E, \Theta \in F_0) + H_{\max}(X|B, \Theta \in F_0) \geq n \log \frac{1}{c}.$$

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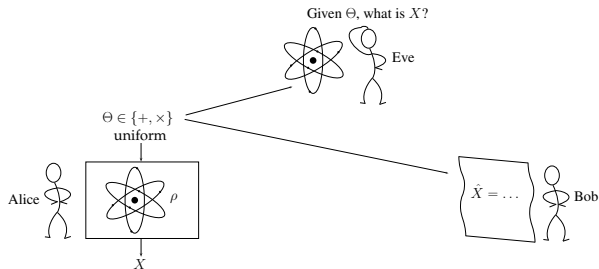
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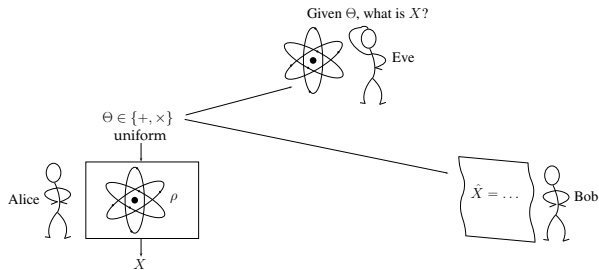
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- Analogously for $F_1 := \{ 'x + + \dots +', '+ x x \dots x' \}$, and so on. Then take the average **and use data-processing** to get

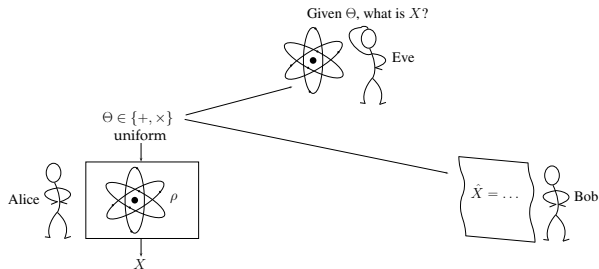
$$H_{\min}(X|E, \Theta) + H_{\max}(X|\hat{X}) \geq n \log \frac{1}{c}.$$

Security Proof Sketch



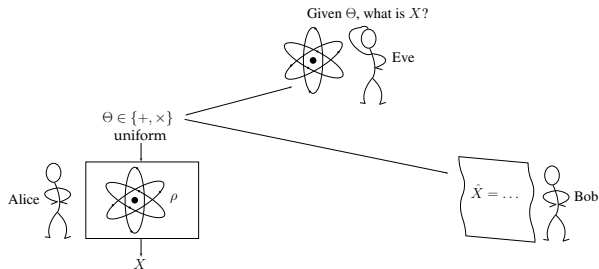
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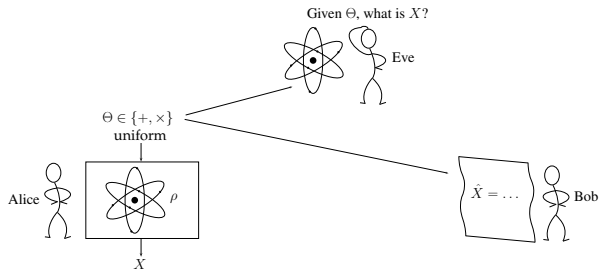
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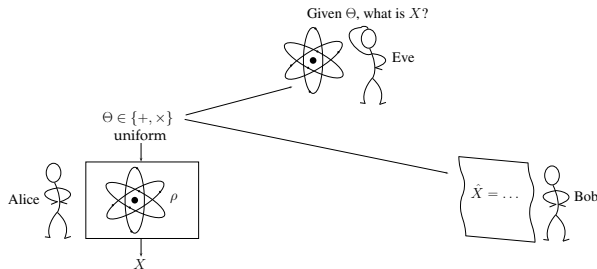
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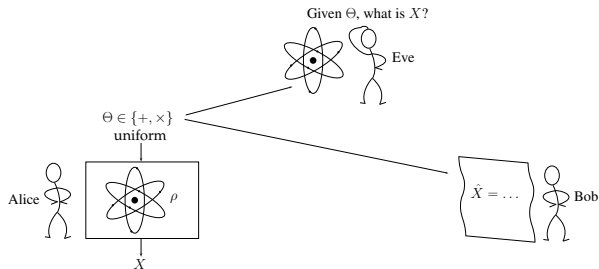
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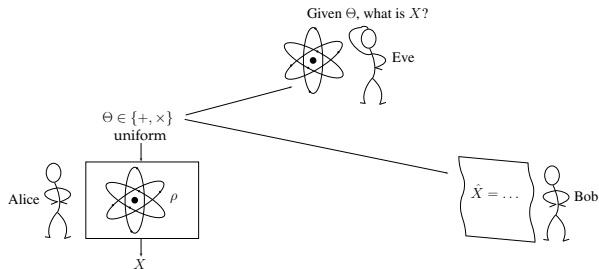
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- Then, estimate $H_{\max}^{\epsilon}(X|\hat{X}) \lesssim nh(\lambda)$.

① Entropic Uncertainty Relations

A Simple Model

Uncertainty for Classical Observers

Maassen & Uffink Relation

Rényi Entropies and the Min-Entropy

② Uncertainty for Quantum Observers

Uncertainty vs. Entanglement

Uncertainty Relation for the Min-Entropy

Quantum Rényi Entropies

Generalized Maassen & Uffink Relation

③ Application to Quantum Key Distribution

Setup for Quantum Key Distribution

Security Proof Sketch

④ More Entropic Uncertainty Relations

Uncertainty Relations for Several Measurements

Various uncertainty relations for classical observers are used to prove security of two-party quantum cryptography, for example in the bounded storage model [DFSS08].

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- See also the survey by Wehner & Winter [WW10].

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- This has applications in continuous variable quantum cryptography (CV-QKD). In Furrer *et al.* [FFB⁺12], we show security for a protocol based on squeezed Gaussian states.
 - In the **finite key length regime** and for **general attacks!**
 - Previous security proofs of CV-QKD only provided security under unrealistic assumptions on the eavesdropper.

Take Home Message

Entropic uncertainty relations allow us to formalize physical intuition in information theoretic terms.

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The min-entropy is the preferred measure of uncertainty for cryptographic applications.

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Uncertainty relations in terms of the min-entropy consequently have a wide range of applications in quantum cryptography.

Bibliography I

- [Ari73] S. Arimoto, *On the Converse to the Coding Theorem for Discrete Memoryless Channels*, IEEE Trans. Inf. Theory **19** (1973), no. 3, 357–359 (English).
- [BB84] Charles H. Bennett and Gilles Brassard, *Quantum Cryptography: Public Key Distribution and Coin Tossing*, Proc. IEEE Int. Conf. Comp., Sys. Signal Process. (Bangalore), IEEE, 1984, pp. 175–179.
- [BBM92] Charles Bennett, Gilles Brassard, and N. Mermin, *Quantum Cryptography Without Bells Theorem*, Phys. Rev. Lett. **68** (1992), no. 5, 557–559.
- [BBR88] Charles H. Bennett, Gilles Brassard, and Jean-Marc Robert, *Privacy Amplification by Public Discussion*, SIAM J. Comput. **17** (1988), no. 2, 210.
- [BCC⁺10] Mario Berta, Matthias Christandl, Roger Colbeck, Joseph M. Renes, and Renato Renner, *The Uncertainty Principle in the Presence of Quantum Memory*, Nat. Phys. **6** (2010), no. 9, 659–662.
- [BCF⁺13] Mario Berta, Matthias Christandl, Fabian Furrer, Volkher B. Scholz, and Marco Tomamichel, *Continuous Variable Entropic Uncertainty Relations in the Presence of Quantum Memory*.
- [CCYZ12] Patrick J. Coles, Roger Colbeck, Li Yu, and Michael Zwolak, *Uncertainty Relations from Simple Entropic Properties*, Phys. Rev. Lett. **108** (2012), no. 21, 210405.
- [Deu83] David Deutsch, *Uncertainty in Quantum Measurements*, Phys. Rev. Lett. **50** (1983), no. 9, 631–633.
- [DFSS08] Ivan B. Damgård, Serge Fehr, Louis Salvail, and Christian Schaffner, *Cryptography in the Bounded-Quantum-Storage Model*, SIAM J. Comput. **37** (2008), no. 6, 1865.
- [FFB⁺12] Fabian Furrer, Torsten Franz, Mario Berta, Anthony Leverrier, Volkher B. Scholz, Marco Tomamichel, and Reinhard F. Werner, *Continuous Variable Quantum Key Distribution: Finite-Key Analysis of Composable Security against Coherent Attacks*, Phys. Rev. Lett. **109** (2012), no. 10, 100502.

Bibliography II

- [Hei27] W. Heisenberg, *Über den anschaulichen Inhalt der quantentheoretischen Kinematik und Mechanik*, Z. Phys. **43** (1927), no. 3-4, 172–198.
- [ILL89] R. Impagliazzo, L. A. Levin, and M. Luby, *Pseudo-random generation from one-way functions*, Proc. ACM STOC, ACM Press, 1989, pp. 12–24.
- [IZ89] R. Impagliazzo and David Zuckerman, *How to Recycle Random Bits*, Proc. IEEE Symp. Found. Comp. Sc., 1989, pp. 248–253.
- [Ken27] E. H. Kennard, *Zur Quantenmechanik einfacher Bewegungstypen*, Z. Phys. **44** (1927), no. 4-5, 326–352.
- [KRS09] Robert König, Renato Renner, and Christian Schaffner, *The Operational Meaning of Min- and Max-Entropy*, IEEE Trans. Inf. Theory **55** (2009), no. 9, 4337–4347.
- [MLDS⁺13] Martin Müller-Lennert, Frédéric Dupuis, Oleg Szehr, Serge Fehr, and Marco Tomamichel, *On quantum Rényi entropies: a new generalization and some properties*.
- [MU88] Hans Maassen and J. Uffink, *Generalized Entropic Uncertainty Relations*, Phys. Rev. Lett. **60** (1988), no. 12, 1103–1106.
- [NBW12] Nelly Huei Ying Ng, Mario Berta, and Stephanie Wehner, *Min-Entropy Uncertainty Relation for Finite-Size Cryptography*, Phys. Rev. A **86** (2012), no. 4, 042315.
- [NJM⁺12] Nelly Huei Ying Ng, Siddarth K Joshi, Chia Chen Ming, Christian Kurtsiefer, and Stephanie Wehner, *Experimental implementation of bit commitment in the noisy-storage model.*, Nat. Commun. **3** (2012), 1326.
- [Rén61] A. Rényi, *On Measures of Information and Entropy*, Proc. Symp. Math., Stat. Probab. (Berkeley), University of California Press, 1961, pp. 547–561.
- [Ren05] Renato Renner, *Security of Quantum Key Distribution*, Ph.D. thesis, ETH Zurich, December 2005.

Bibliography III

- [Sha48] C. Shannon, *A Mathematical Theory of Communication*, Bell Syst. Tech. J. **27** (1948), 379–423.
- [TCR09] Marco Tomamichel, Roger Colbeck, and Renato Renner, *A Fully Quantum Asymptotic Equipartition Property*, IEEE Trans. Inf. Theory **55** (2009), no. 12, 5840–5847.
- [TR11] Marco Tomamichel and Renato Renner, *Uncertainty Relation for Smooth Entropies*, Phys. Rev. Lett. **106** (2011), no. 11.
- [WW08] Stephanie Wehner and Andreas Winter, *Higher entropic uncertainty relations for anti-commuting observables*, J. Math. Phys. **49** (2008), no. 6, 062105 (en).
- [WW10] ———, *Entropic Uncertainty Relations—A Survey*, New J. Phys. **12** (2010), no. 2, 025009.
- [WWY13] Mark M. Wilde, Andreas Winter, and Dong Yang, *Strong Converse for the Classical Capacity of Entanglement-Breaking and Hadamard Channels*.